

- 1) There will be an induced surface charge density on the surface of the sphere that produces exactly the same potential outside the sphere as a mirror charge. Thus outside the sphere the potential is that from the charge q and the mirror charge. The mirror charge has the charge $-aq/b$ and is placed at the distance a^2/b from the origin. Do we have any monopole contribution? Yes, there is a net charge and the monopole contribution is

$\Phi^{(1)}(r, \theta, \varphi) = [q - (aq/b)]/r = q[1 - (a/b)]/r$. Do we have any dipole contribution? Yes the dipolemoment is:

$$\mathbf{p} = \hat{z} [qb - (aq/b)(a^2/b)] = \hat{z} qb [1 - (a^3/b^3)] \text{ and } \Phi^{(2)}(r, \theta, \varphi) = \mathbf{p} \cdot \mathbf{r} / r^3 = pr \cos \theta / r^3 = qb [1 - (a^3/b^3)] \cos \theta / r^2$$

Do we have axial symmetry? Yes, and

$$Q = q(3b^2 - b^2) - (aq/b) [3(a^2/b)^2 - (a^2/b)^2] = q[2b^2 - 2a^5/b^3] = 2qb^2 [1 - (a/b)^5]; \Phi^{(4)}(r, \theta, \varphi) = Q(3 \cos^2 \theta - 1) / 4r^3$$

$$\Phi(r, \theta, \varphi) = q[1 - (a/b)]/r + qb [1 - (a^3/b^3)] \cos \theta / r^2 + 2qb^2 [1 - (a/b)^5] (3 \cos^2 \theta - 1) / 4r^3$$

2. The potential from a line charge is obtained from Gauss' law:

$$2\pi r \Delta l D(r) = 4\pi \sigma_l \Delta l \rightarrow E(r) = D(r) = 2\sigma_l \Delta l / (r \Delta l) = \sigma_l / r \rightarrow \Phi(r) = -2\sigma_l \ln(r) = -\sigma_l \ln(r^2)$$

We are supposed to show that the total potential outside the cylinder can be written as:

$$\begin{aligned} \Phi(\mathbf{r}) &= -\sigma_l \ln [x^2 + (y-b)^2] + \sigma_l \ln [x^2 + (y - a^2/b)^2] + C_1 \\ &= -\sigma_l \ln \left\{ \frac{[x^2 + (y-b)^2]}{[x^2 + (y - a^2/b)^2]} \right\} + C_1 = -\sigma_l \ln \left\{ C_2 \frac{[x^2 + (y-b)^2]}{[x^2 + (y - a^2/b)^2]} \right\} \end{aligned}$$

We should show that the proper choice of constant leads to the potential zero on the whole cylinder surface.

$$\begin{aligned} \Phi(r=a) &= |x^2 = a^2 - b^2| = -\sigma_l \ln \left\{ C_2 \frac{[a^2 - y^2 + (y-b)^2]}{[a^2 - y^2 + (y - a^2/b)^2]} \right\} \\ &= -\sigma_l \ln \left\{ C_2 \frac{[a^2 + b^2 - 2yb]}{[a^2 + a^4/b^2 - 2ya^2/b]} \right\} \\ &= -\sigma_l \ln \left\{ C_2 \frac{[a^2 + b^2 - 2yb]}{b^2} \frac{[b^2 + a^2 - 2yb]}{b^2} \right\} = -\sigma_l \ln \left\{ C_2 \frac{b^2}{a^2} \right\} = 0 \rightarrow C_2 = \frac{a^2}{b^2} \rightarrow C_1 = \sigma_l \ln \left(\frac{b^2}{a^2} \right) \end{aligned}$$

- 3.

$$\Phi(u, v) = -V + \frac{2V}{\pi} v = -V + \frac{2V}{\pi} \arg(z) = -V + \frac{2V}{\pi} \tan^{-1}(y/x) \rightarrow \mathbf{E} = -\nabla \Phi = -\frac{2V}{\pi} \frac{1}{1 + (y/x)^2} (-y/x^2, 1/x)$$

$$\rightarrow E_y(x, 0^+) = \frac{-2V}{\pi x} \rightarrow \sigma = \frac{-2V}{\pi x 4\pi} = \frac{-V}{2\pi^2 x}$$

4. The velocity of the charges (drift velocity due to the rotation) is $\mathbf{v} = \omega R \hat{\phi}$. They are influenced by a Lorentz force $\mathbf{F} = q\mathbf{v} \times \mathbf{B}/c = q\omega R \hat{\phi} \times B \hat{z}/c = q\omega R B \hat{\phi} \times \hat{z}/c = q\omega R B \hat{r}/c$. An EMF is developed from the center towards the rim: $\mathcal{E} = \frac{1}{c} \int_0^a \omega R B dR = \omega B a^2 / 2c$. In other words there will be a potential difference of this size across the resistor and the resulting current is: $I = \omega B a^2 / 2Rc$. If we do not have a resistor a large current can be generated. It is limited by the total resistance in the circuit. The power loss in the resistor is $RI^2 = I \omega B a^2 / 2c$ which is equal to the supplied power, i.e., the torque times the angular velocity: $N\omega = I \omega B a^2 / 2c \rightarrow N = I B a^2 / 2c$

Here we can imagine some applications. If the \mathbf{B} -field is known we may have an ampere-meter displaying torque. Alternatively one may measure magnetic induction if one knows I and N .